

**Exercise 1: The metric tensor in spherical coordinates**

**(6 Points)**

Consider a coordinate transformation in Minkowski space from cartesian to spherical space coordinates,  $x^\mu = (t, x, y, z)^T \rightarrow \xi^\mu = (t, r, \theta, \varphi)^T$ . The line element in the first set of coordinates reads  $ds^2 = -dt^2 + dx^2 + dy^2 + dz^2$ .

- Compute the metric tensor,  $g_{\mu\nu}$ , in spherical space coordinates. Write down the line element in the new coordinates.
- Compute the inverse metric tensor,  $g^{\mu\nu}$ , in both sets of coordinates.
- Consider now Minkowski space restricted to the surface of a sphere with radius  $R$ . Use the previous results to compute the metric on this surface. Can you find a coordinate transformation that would convert this metric into its cartesian form,  $ds^2|_{r=R} = -dt^2 + dx^2 + dy^2$ ?

Make your calculations explicit.

**Exercise 2: The action of SR**

**(7 Points)**

The action of a free particle is minimized by the path with the shortest distance between two points. Since distances for timelike particles are measured by the proper time the action of a free particle in special relativity can be expressed in the following way:

$$S = \alpha \int d\tau = \int L dt, \quad (1)$$

where the proper time  $\tau$  is defined by

$$d\tau^2 = -\eta_{\mu\nu} dx^\mu dx^\nu. \quad (2)$$

- Determine the relation between the proper time interval  $d\tau$  and the time interval  $dt$  by using eq. (2).
- Identify  $\alpha$  by calculating the non-relativistic limit  $v \ll 1$ .
- Use the Euler-Lagrange equation to find the equation of motion for the relativistic case.

We now consider the general case, which yields the following action:

$$S = \alpha \int \sqrt{-\eta_{\mu\nu} dx^\mu dx^\nu} = \alpha \int \sqrt{-\eta_{\mu\nu} \frac{dx^\mu}{d\tau} \frac{dx^\nu}{d\tau}} d\tau. \quad (3)$$

Here,  $x^\mu = x^\mu(\tau)$ .

- Use the Euler-Lagrange equation again for the Lagrangian  $L(x^\mu, \frac{dx^\mu}{d\tau}, \tau)$  in eq. (3) to obtain the equations of motion

$$\frac{d^2 x^\mu}{d\tau^2} = 0. \quad (4)$$

- e) Finally, show that the equations of motion in a different coordinate system with coordinates  $y^\mu$  are of the form

$$\frac{d^2 y^\mu}{d\tau^2} = - \frac{\partial y^\mu}{\partial x^\sigma} \frac{\partial^2 x^\sigma}{\partial y^\nu \partial y^\rho} \frac{dy^\nu}{d\tau} \frac{dy^\rho}{d\tau}. \quad (5)$$

This equation is called the *geodesic equation*.

**Exercise 3: Perfect Fluid**

**(7 Points)**

- (a) Explain the term perfect fluid. What are the properties of the energy-momentum tensor  $T^{\mu\nu}$  for a perfect fluid?
- (b) Which important conservation law does the energy-momentum tensor satisfy? Calculate  $\partial_\mu T^{\mu\nu}$  explicitly for a perfect fluid.
- (c) Consider the projection tensor  $P^\sigma{}_\nu = \delta^\sigma{}_\nu + u^\sigma u_\nu$ . This tensor projects a vector onto a vector which is orthogonal to the fluid's four-velocity. Check whether  $P^\sigma{}_\nu$  indeed is the projection tensor by calculating  $P^\sigma{}_\nu P^\nu{}_\mu$  and  $P^\sigma{}_\nu u^\nu$ .
- (d) Project the vector obtained in (b) onto a vector which is orthogonal to the fluid's four-velocity by using the projection tensor. What familiar equation from classical fluid mechanics do you find in the non-relativistic limit?