**<http://people.het.physik.tu-dortmund.de/~ghiller/WS1920ETT.html>**

## **Exercise 1: Adjoint representation of** *SU*(*N*) **(8 Points)**

The generators  $T^a$  of the group  $SU(N)$  are hermitian operators which construct all infinitesimal group transformations. The commutator of two generators can be written as a linear combination

$$
\left[T^a, T^b\right] = \mathrm{i} f^{abc} T^c \,,\tag{1.1}
$$

where  $f^{abc}$  are structure constants. The commutator relation [\(1.1\)](#page-0-0) and the vector space spanned by the generators are called the *Lie algebra* of the corresponding Lie group *SU*(*N*).

(a) Derive the *Jacobi identity*

$$
\left[T^{a},\left[T^{b},T^{c}\right]\right]+\left[T^{b},\left[T^{c},T^{a}\right]\right]+\left[T^{c},\left[T^{a},T^{b}\right]\right]=0
$$
\n(1.2)

using the following relation

$$
f^{ade} f^{bcd} + f^{bde} f^{cad} + f^{cde} f^{abd} = 0.
$$
 (1.3)

(b) For every simple Lie algebra an adjoint representation can be defined. The corresponding generators can be written as

$$
\left(t^b\right)_{ac} = \mathrm{i} f^{abc} \,. \tag{1.4}
$$

Show that these generators satisfy the Lie algebra

$$
\left(\left[t^{a}, t^{c}\right]\right)_{be} = \mathrm{i} f^{acd} \left(t^{d}\right)_{be}.\tag{1.5}
$$

Use Eq. [\(1.3\)](#page-0-1) and keep in mind that the structure constants are anti-symmetric in the first two indices, i.e.  $f^{abc} = -f^{bac}$ .

(c) The quadratic Casimir invariant of a representation *R* is defined as

$$
C(R)\mathbb{1} = \sum_{a} \left(t^{a} t^{a}\right),\tag{1.6}
$$

where  $t^A$  denotes the generators of the gauge group in a representation R. The Dynkin index of a representation *R* reads

$$
S(R)\,\delta^{ab} = \text{Tr}\left(t^a\,t^b\right). \tag{1.7}
$$

Show that for the adjoint representation of the *SU*(3) the following relation holds

$$
C(\text{adj}) = S(\text{adj}) = 3,\tag{1.8}
$$

taking into consideration the non-vanishing *SU*(3) structure constants

$$
f^{123} = 1, f^{147} = -f^{156} = f^{246} = f^{257} = f^{345} = -f^{367} = \frac{1}{2}, f^{458} = f^{678} = \frac{\sqrt{3}}{2}.
$$
 (1.9)

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## **Exercise 2: Parton distribution functions (6 Points)**

- a) On the website http://hepdata.cedar.ac.uk/pdf/pdf3.html you can find the parton distribution functions  $xf(x,Q^2)$  as a function of the Bjorken variable *x*. Use the data of the (G)JR group, set JR09FFnnlo, and plot the parton distribution functions of the quarks and gluons with  $Q^2 \in \{0.1, 1, 10\}$  GeV<sup>2</sup>. Interpret the plots.
- b) Sum rules can be derived for the parton distribution functions. From the definition of the total momentum of the proton one finds

$$
\int_0^1 x \sum_i f_i(x) = 1.
$$
 (2.1)

Consider the distribution functions  $u<sub>v</sub>(x)$ ,  $d<sub>v</sub>(x)$ ,  $s(x)$ ,  $\bar{s}(x)$  and  $g(x)$ , where the index *v* denotes the valence quarks and non-indexed quarks are virtual sea quarks. What are the sum rules of the electric charge  $Q = +1$  and the vanishing strangeness *S* = 0 of the proton?

## **Exercise 3: Isospin rotations for quarks and antiquarks (6 Points)**

The up- and down-quark fields can be written as SU(2)-Isospin doublets  $q = (u, d)^T$ , where for a general isospin doublet  $(q_1, q_2)^T$  the components  $q_1$  and  $q_2$  fulfil  $I_3(q_1) = 1/2$ and  $I_3(q_2) = -1/2$ . Here  $I_3$  denotes the third component of the isospin. In the following we want to investigate the transformation of such a doublet under a rotation*U* in isospin space. The rotation is given by

<span id="page-1-0"></span>
$$
q' = Uq \quad \text{with} \quad U = \exp(i\vec{\phi}\vec{\tau}), \tag{3.1}
$$

with angles  $\phi_i$  and rotation generators  $\tau_i = \sigma_i/2$ , where  $\sigma_i$  are the Pauli matrices.

- a) Calculate *U* and  $q' = Uq$  for a rotation around  $\phi = (0, \pi, 0)^T$ .
- b) Consider now the antiquark doublet. The antiquarks belong to the so-called conjugate representation of SU(2), so that for the isospin part of the fields the relations  $\overline{u} = u^*$  and  $\overline{d} = d^*$  apply (the quotation marks indicate that the relations are not valid for the entire fields, but only for their respective isospin part). Because of the Gell-Mann-Nishijima relation,  $I_3(\bar{u}) = -I_3(u)$  and  $I_3(\bar{d}) = -I_3(d)$  hold as well. Therefore, the ansatz  $q = (a\bar{d}, b\bar{u})^T$  may be used for complex numbers  $a, b$  with  $|a| = |b| = 1$ . Proceed as follows: Conjugate equation [\(3.1\)](#page-1-0) and express the result using this ansatz for the antiquark doublets. How should you choose *a* and *b* so that the antiquark doublet is actually a doublet, that is,  $\bar{q}' = U\bar{q}$  is fulfilled?

Note: Use your result from subtask a) to determine the coefficients.